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Latitudinal extent of large-scale structures in the solar wind

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Abstract. Comparison of solar wind observations from the ACE spacecraft, in the ecliptic plane at ~ 1 AU, and the Ulysses spacecraft as it orbits over the Sun's poles, provides valuable information about the latitudinal extent and variation of solar wind structures in the heliosphere. While qualitative comparisons can be made using average properties observed at these two locations, the comparison of specific, individual structures requires a procedure to determine if a given structure has been observed by both spacecraft. We use a 1-D hydrodynamic code to propagate ACE plasma measurements out to the distance of Ulysses and adjust for the differing longitudes of the ACE and Ulysses spacecraft. In addition to comparing the plasma parameters and their characteristic profiles, we examine suprathermal electron measurements and magnetic field polarity to help determine if the same features are encountered at both ACE and Ulysses. The He I λ 1083 nm coronal hole maps are examined to understand the global structure of the Sun during the time of our heliospheric measurements. We find that the same features are frequently observed when both spacecraft are near the ecliptic plane. Stream structures derived from smaller coronal holes during the rising phase of solar cycle 23 persists over 20° – 30° in heliolatitude, consistent with their spatial scales back at the Sun.

Key words. Interplanetary physics (solar wind plasma)

1 Introduction

Coronagraph measurements provide information about the size and three-dimensional shape of coronal structures within $\sim 30 R_S$ of the Sun, but further out in the solar wind, much less is known about how such large-scale structures evolve. Over the past forty years, most of what has been learned about the heliosphere has been based on in-ecliptic measurements. In contrast, the unique orbit of Ulysses, which passes over the Sun's poles, provides in situ measurements as a function of latitude. Ulysses has already greatly enhanced

our knowledge of the 3-D structure of the inner heliosphere during solar minimum when the large-scale structure varies slowly (e.g. McComas et al., 2000, and references therein). For example, Ulysses solar minimum measurements revealed a relatively simple heliospheric structure with high-speed streams emitted from large polar coronal holes filling the high latitude heliosphere. In contrast, near-maximum solar wind is much more complex with highly variable flows arising from a variety of coronal sources observed at all heliolatitudes (McComas et al., 2002a).

In this paper we have two goals: (1) to develop and test techniques for determining when two spacecraft observe the same large-scale structures, such as coronal hole flows and Coronal Mass Ejections (CMEs) in the heliosphere, and (2) to determine the latitudinal extent of these structures in the heliosphere.

Hydrodynamic codes have been used to study the solar wind for a long time, and simulations have advanced to the point of being 3-dimensional magnetohydrodynamic codes with solar observations supplying inner boundary conditions (e.g. Riley et al., 2002). We use 1-D hydrodynamic simulations driven at the inner boundary with spacecraft observations, and adjust the data in time to account for corotation with the Sun. In some Helios (e.g. Schwenn et al., 1981) and Pioneer (e.g. Mitchell et al., 1981) studies, observations from pairs of spacecraft were rotated to align measurements in longitude, but they used a constant solar wind speed for propagation between spacecraft. Very few studies have used a 1-D code with a driven inner boundary. We know of only one other study by De Keyser et al. (2000) that uses such a technique to examine sector boundaries. However, these authors limited their study to the rare times when WIND (1 AU) and Ulysses (5 AU) were in radial alignment. Our study is unique in that we use Ulysses and ACE data when the spacecraft have latitude separations from 0 – 42° , and then use the sector information as a tool to examine the latitude extent of large-scale structures, and their dynamic evolution.

Several other recent studies used different approaches to examine large-scale features. Recently, Neugebauer et al.

(2002) compared Ulysses and ACE data to solar observations to determine the source site on the Sun of the different types of solar wind. We use both ACE and Ulysses data as do Neugebauer et al. (2002), but we focus on the dynamics of large-scale structures. Recent work by Riley et al. (2002) provides motivation for examining the dynamics of large-scale features in the inner heliosphere. They found that using constant speed mapping produces a significantly different heliospheric current sheet shape than when dynamics are included. Neugebauer et al. (1998) showed that a magnetohydrodynamic (MHD) model and several source-surface models predict boundary positions that differ by as much as 20° . Wang et al. (2000a,b), and Richardson et al. (2002) examined the evolution of structures in the distant heliosphere by comparing Ulysses and Voyager data. We focus our attention on the evolution of structures from 1–5 AU, which is a region where stream interactions grow and the coalescence of structures begins to occur. The slowdown of the solar wind by interstellar neutrals is negligible within 5 AU. The region from 1–5 AU is ideal for studying stream evolution and the development of corotating interaction regions

2 Coronal hole plasma and CMEs

High-speed streams have a characteristic speed profile due to compression and rarefaction produced as the streams interact with surrounding slower plasma. Plasma emanating from coronal holes has other interplanetary signatures in addition to having high speed. These signatures include the proton to alpha ratio, magnetic field structure, heavy ion composition, and ion charge state abundance. High-speed streams tend to consist of one magnetic polarity. This supports the predominant view that high-speed streams ($>700 \text{ km s}^{-1}$) generally come from coronal holes, which also consist of one dominant polarity (Hundhausen, 1977). High-speed streams have low oxygen and carbon freezing-in temperatures and low Mg/O and Fe/O ratios (Geiss et al., 1995), low density (Belcher and Davis, Jr., 1971), relatively structure free plasma parameters (Feldman et al., 1996, and references therein), and Alfvén waves are often present (Belcher and Davis, Jr., 1971). High-speed streams associated with the large polar coronal holes that Ulysses observed around solar minimum, have a mean speed of $\sim 760 \text{ km s}^{-1}$, mean 1 AU temperature of $2.7 \times 10^5 \text{ K}$, mean 1 AU density of 2.7 cm^{-3} , and a mean alpha to proton ratio of 4.4% (McComas et al., 2000).

Here, we present observations during the rising phase of the solar cycle. Around solar maximum large-scale transient coronal mass ejections are more frequent (Webb and Howard, 1994). While there are many in situ signatures of Interplanetary Coronal Mass Ejections (ICMEs), no one signature can be used to identify all ICMEs (Gosling, 1996; Neugebauer and Goldstein, 1997). Recently, these signatures have been reviewed by Gosling (1996, 1997), and Neugebauer and Goldstein (1997). At 1 AU the presence of counterstreaming suprathermal electrons is a fairly reliable

signature of ICMEs (Gosling, 1996). A few other criteria are enhanced helium abundances ($\text{He}^{++}/\text{H}^+ > 0.08$), low beta, strong magnetic fields, and unusual ionization states such as Fe^{+16} .

In addition to changes in the frequency of CMEs with the solar cycle, the number and size of coronal holes vary over the course of the solar cycle. At solar minimum there are two large polar coronal holes and an equatorial streamer belt. However, during the rising phase and near solar maximum the polar coronal holes shrink in size and eventually disappear, and many small coronal holes develop (Wang et al., 1996). These small coronal holes are not limited to high latitudes and occur at mid and low latitudes. Many of these small holes are associated with active regions and transients (Wang et al., 1996). During the descending phase of the solar cycle, polar coronal holes reform, and equatorial extensions to polar coronal holes are more common (Hundhausen et al., 1981). The lifetimes of coronal holes are greater during the descending phase (Hundhausen et al., 1981). These smaller coronal holes produce streams in the heliosphere with speeds typically reaching only $500\text{--}600 \text{ km s}^{-1}$, which is less than the large coronal holes (McComas et al., 2002b). However, McComas et al. (2002b) found the freezing-in temperatures to be characteristic of the larger polar coronal holes. The small coronal hole freezing-in carbon and oxygen temperatures were less than 1.1 MK and 1.3 MK, respectively.

3 Data and model

Primarily, we use ion and electron measurements from the Advanced Composition Explorer (ACE) and Ulysses spacecraft. These observations are taken from the ACE Solar Wind Electron Proton Alpha Monitor (SWEPAM) (McComas et al., 1998) and the Ulysses Solar Wind Observations Over the Poles of the Sun (SWOOPS) instruments (Bame et al., 1992). Supporting observations from the magnetometers on Ulysses (Balogh et al., 1992) and ACE (Smith et al., 1998) are also used. During the time of this study, 5 February 1998 to 31 December 1999, the ACE spacecraft was near Earth at L1, and Ulysses started at 2° latitude and 5.39 AU and went south, ending at -42° latitude and 4.16 AU.

To propagate data from 1 AU (ACE) out to the location of Ulysses, we use the Zeus astrophysical single-fluid magnetohydrodynamic model (Stone and Norman, 1992). The model uses a Eulerian finite difference scheme. Although the Zeus model can include magnetic fields and radiation transport, in this study they are neglected and only hydrodynamic effects are examined. Not including magnetic fields causes waves to travel at the sound speed instead of the magnetosonic speed. A further consequence of not including hydromagnetic effects is that forces due to magnetic pressure gradients and tension are not included (Riley et al., 1997). Since we examine data at radial distances less than 5.4 AU, pickup ions effects are not included. Inclusion of pickup ions would slightly modify the pressure profile (Riley and

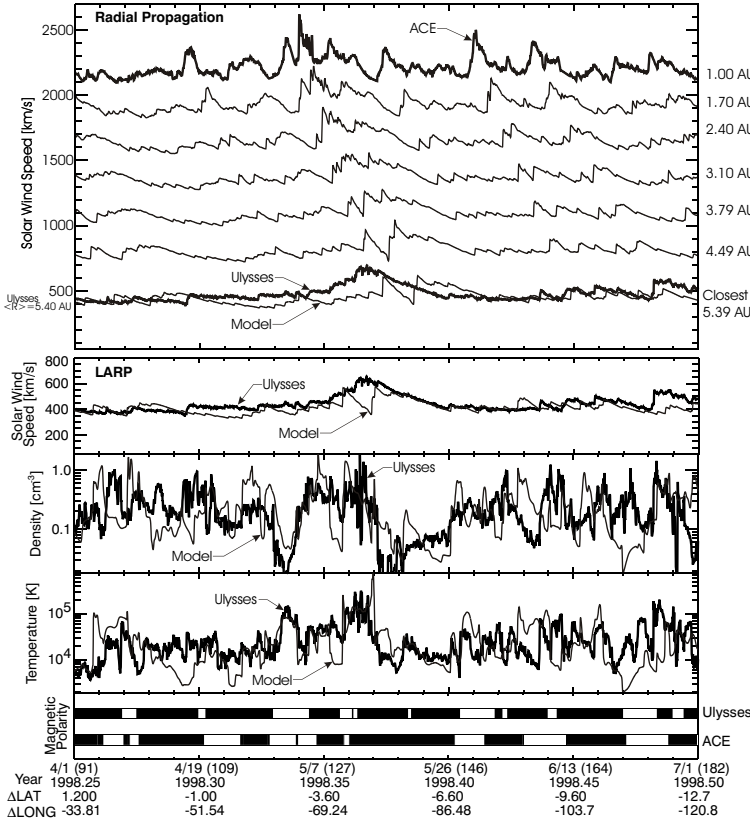


Fig. 1. The top panel shows ACE solar wind speed data at 1 AU as a thick line, and model speeds at given propagation distances as thin lines. The Ulysses data overlays the ACE propagated data at the grid step that most closely matches the distance of Ulysses. Subsequent panels compare Ulysses data and model results when the longitudes of the two spacecraft have been aligned in addition to the radial propagation. The second panel displays the velocity, the third panel the density, the fourth panel the temperature, and the last panel the polarity of the magnetic field. Additional x -axis labels indicate the difference in latitude (ΔLAT) and longitude (ΔLon) between the ACE and Ulysses spacecraft.

Gosling, 1998). The model solves a system of continuity, momentum, and internal energy equations. A polytropic index of $3/2$ is used since most solar wind measurements are generally consistent with this value (Riley et al., 2001). Even though Zeus is a 2-D model, only one dimension is used in this study so that we can use single point in situ measurements to drive the model. Since we know the inner boundary conditions (1 AU) for a point, we use a 1-D version of the model to propagate measurements from ACE at 1 AU to the position of Ulysses when it was at distances between 4.2 and 5.4 AU. We focus on the dynamic evolution of structures between 1 and 5.4 AU without making assumptions about the solar wind properties at locations away from our measurements. Gosling et al. (1995, 1998) showed that a 1-D model predicts many key features of ICMEs, which are also found in more complex 2-D models of ICMEs (Riley et al., 1997). However, a 1-D model cannot accurately predict complex effects like those observed by Riley et al. (1997) in their 2-D model. Riley et al. (1997) found that latitudinal gradients in the ambient solar wind speed can affect the evolution of an ICME to the extent that an ICME may split in two. A 1-D model would not be able to predict such a split. De Keyser et al. (2000) found that by using a 1-D model they could often identify sector boundaries observed at Ulysses in the WIND data when Ulysses and WIND were in radial alignment.

We use hourly averages of the Ulysses-SWOOPS and ACE-SWEPAM plasma measurements. The ACE measurements are used as the inner boundary conditions for each time

step of the simulation and then we compare the propagated ACE measurements with the Ulysses measurements. Since the Ulysses spacecraft position varies slowly around apheion, the data are split into 4 segments per year to calculate an average radial position. For each segment we compare Ulysses SWOOPS measurements to ACE mapped data at the radial grid point that most closely matches the average radial distance of Ulysses. The grid size is $0.47 R_S$. We use the magnetic polarity, magnitudes and profiles of the velocity, density, and temperature to identify specific features. The polarity is determined by comparing the magnetic field direction to the Parker spiral direction determined using the solar wind speed (Forsyth et al., 1996). Since we are interested in the large-scale structure, a 24-hour running mean is used to calculate the sector structure. We also calculated a 12 h running mean and obtained very similar results. We did not use a short time period for the averaging because then the Alfvén waves would produce polarity changes. Suprathermal electron distributions are used to identify coronal mass ejections.

4 Results

The top panel of Fig. 1 displays a time series of ACE-SWEPAM solar wind speed measurements at 1 AU (top [4] curve), the radial propagation (RP) of the ACE data at intermediate grid steps (5 middle curves), and a direct comparison between Ulysses SWOOPS measurements and mapped ACE

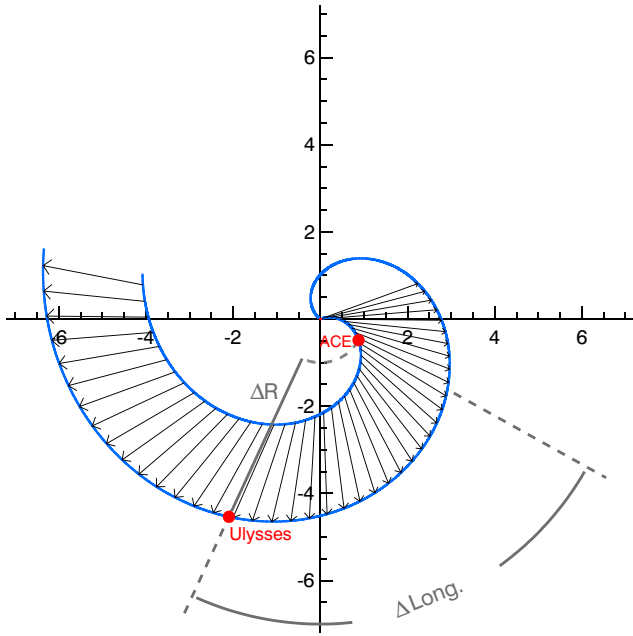


Fig. 2. Diagram of LARP method. Blue lines are magnetic field lines, red dots are spacecraft, and black arrows are velocity vectors. The Sun is at the center and the units of distance are AU.

data (bottom curve). Each of the successive curves is offset by 40 km s^{-1} in this stacked format. The compressions on the leading edges of high-speed streams become steeper with distance as faster plasma overtakes slower plasma ahead. Such compressions on the leading edges of corotating interaction regions typically steepen into shocks between 1.5 and 2.5 AU, and can easily be tracked out to distances greater than 5 AU. Through this process, high-speed structures are worn down as momentum and energy are transferred to lower speed plasma. For example, from 1998.24 to 1998.34 and from 1998.4 to 1998.47 the model predicts that the high-speed features are worn down, consistent with Ulysses measurements of a more uniform low speed solar wind speed. Between 1998.34 and 1998.4, Ulysses measurements show enhanced speeds, but the peak in the mapped data occurs at a later time. For this large feature between 1998.34–1998.4, ACE and Ulysses electron measurements both show the presence of counterstreaming electrons, indicating that this structure is an ICME. Gosling (1996) reviewed signatures of ICMEs in the heliosphere and found that the presence of counterstreaming electrons is perhaps the most reliable signature of an ICME.

The RP method does not take into account the longitudinal separation of the two spacecraft, which varies between 0° and 360° each year, as the Earth and L1 revolve around the Sun. To compensate for this effect (at least for co-rotating structures), one set of observations has been “de-rotated” with respect to the other. We will refer to the combination of longitude adjustment and radial propagation as LARP. The propagation time for the LARP method can be thought of as the time required for a spiral at ACE to propagate to Ulysses,

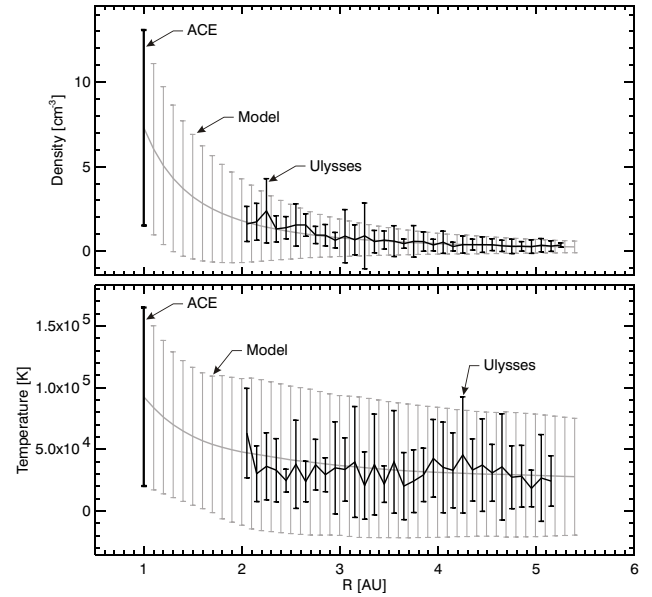


Fig. 3. Density and temperature data divided into radial bins and averaged. Simulation results are shown in grey, ACE and Ulysses data are shown in black. Description of the data set and binning is provided in the text.

as shown schematically in Fig. 2. To examine Ulysses observations at higher latitudes, the LARP method may need to be refined to include the latitude dependence of the spiral winding. The second panel in Fig. 1 overlays the speed profiles after shifting the radially propagated ACE data in time, in order to match the ACE source longitudes to the Ulysses ones. Practically, this was accomplished by subtracting the inertial heliospheric longitudes for the two spacecraft and converting this difference into an effective time shift using the solar rotation rate. In order to minimize the effects of temporal evolution of the solar source, the data are rotated the shortest way around (forward or backward in time) so that comparisons are always made between data separated by less than two weeks. Using the LARP method causes the large feature in the Ulysses data at 1998.34 to align with the double peaked feature in the mapped data.

It appears that the series of peaks observed at ACE either do not merge properly in the model, or parts of an ICME could evolve differently as some ICME model results show (Riley et al., 1997). The LARP method works particularly well for the ICME at 1998.34. In total, we have examined 15 counterstreaming intervals from 1 February 1998 to 31 October 1998, and find that the LARP method improves the sector alignment for 10 of those intervals. There were not enough CME observations to draw definitive conclusions. If an ICME stays magnetically connected to the Sun for a long time, then it becomes aligned along a spiral, as depicted by McComas et al. (1992), such that the westward flank becomes elongated more than the eastward flank. Also, the spiral angle of such an ICME, as with any magnetically connected solar wind flow, depends on the speed of the ICME.

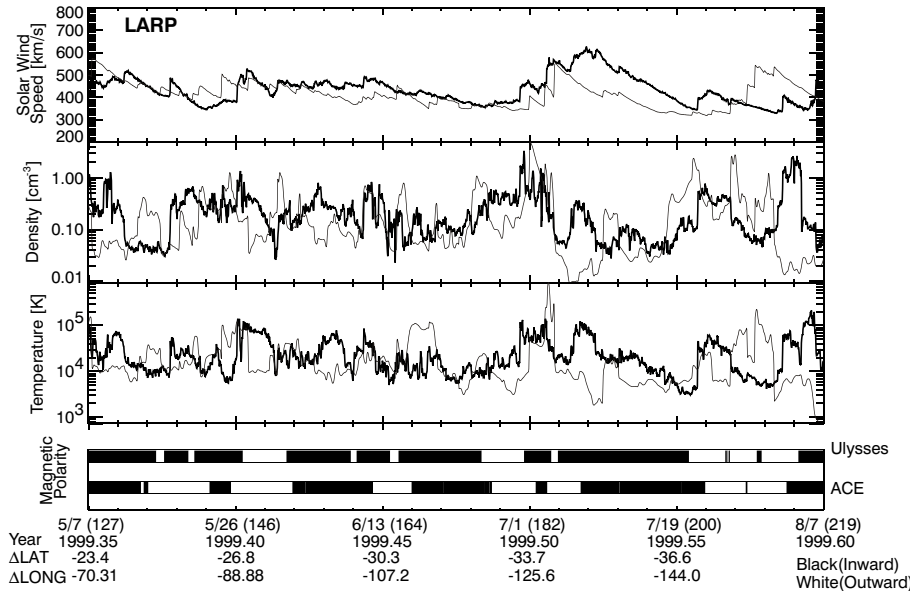


Fig. 4. Similar format as Fig. 1, but is at a time where ACE and Ulysses have a larger separation in latitude. The top panel shows the solar wind speed ACE data propagated using the LARP method (thin black line) and the Ulysses data (thick black line). Density is shown in the second panel, temperature in the third and sector in the fourth panel.

The bottom three panels of Fig. 1 compare Ulysses density (third panel), temperature (fourth panel), and magnetic polarity (bottom panel) measurements to the LARP ACE data. The LARP ACE speeds, average density, and temperature agree fairly well with the Ulysses data; however, the high frequency fluctuations in density and temperature do not. We would not expect high frequency fluctuations to agree well because while the two spacecraft observe the same large-scale solar wind structures, they never observe precisely the same parcels of plasma. ACE and Ulysses show very similar sector structure throughout this entire interval. From 1998.31 to 1998.42 the mean mapped density and variations in density agree well with the data. The mapped mean temperature and temperature profile agrees well with the measured temperature over a longer time period from 1998.31 to 1998.5, although from 1998.42 to 1998.5 the difference between the model and data average temperature is greater than earlier. It appears that the mapped density and temperatures agree less well than the velocity and sector mainly because the temperature and density are such highly variable parameters.

In order to compare the average variations in the mapped density and temperature with the Ulysses observations, we calculated radial profiles. In Fig. 3 the model results from 5 February 1998 through 31 December 1999 are binned in radial distance and averaged over time. The ACE and Ulysses measurements are analyzed similarly, except Ulysses data are taken from 31 December 1997 to 3 March 2001, when Ulysses was moving over radial distances ranging from 1.6 to 5.4 AU. The radial variation of density and temperature observed in the Ulysses data are consistent with the density being proportional to R^{-2} and temperature proportional to R^{-1} , as shown in a statistical study of Ulysses' entire first orbit observations (McComas et al., 2000) and calculated in this study by the Zeus model. In the model the R^{-2} dependence is due to spherical expansion, and the temperature de-

pendence is then R^{-1} for a polytropic index of 3/2.

As the latitude separation between the spacecraft becomes larger, ACE and Ulysses do not observe the same features as frequently. The sector structure at ACE and Ulysses are at times similar, however, the plasma parameters do not agree over such large time spans as when the latitude separation was smaller. For example, Fig. 4 shows the interval from 1999.35 to 1999.65, as the latitude separation between the two spacecraft grew from 20° to 38° . The model speed agrees well at the beginning (1999.44–1999.51) when the speed is steady, but the general shape and extent of the first peak (1999.51) does not match well. This may indicate that our mapping has correctly aligned the structures in longitude, and the differences reflect the spatial differences of the structure, that is, ACE and Ulysses are probing different parts of a given CIR. Different parts of a CIR or ICME probably have different compression strengths and thicknesses. Towards the end of this plot the model predicts a peak in velocity at 1999.58 that is not observed at Ulysses. It appears that when ACE and Ulysses have large latitude separations the two spacecraft do not encounter the same features as frequently as they do at smaller latitude separations.

Inspection of the coronal hole maps shown in Fig. 5, and the electron distributions leads us to believe that we may not be observing the same structures in 1999 when the spacecraft become separated by more than 27° in latitude. The coronal hole maps in Fig. 5 are derived from He I $\lambda 1083$ nm Kitt Peak heliospectrograms. Harvey (1996) and Harvey and Recely (2002) use four criteria to determine the boundaries. Those criteria are the coronal holes (1) appear bright in the He I $\lambda 1083$ nm spectroheliograms, (2) contain between 75% and 100% of one magnetic polarity, (3) are at least 2 supergranules in size, and (4) have low network contrast. Coronal hole boundaries expand in latitude and longitude with altitude in the corona, but we still find these maps to provide

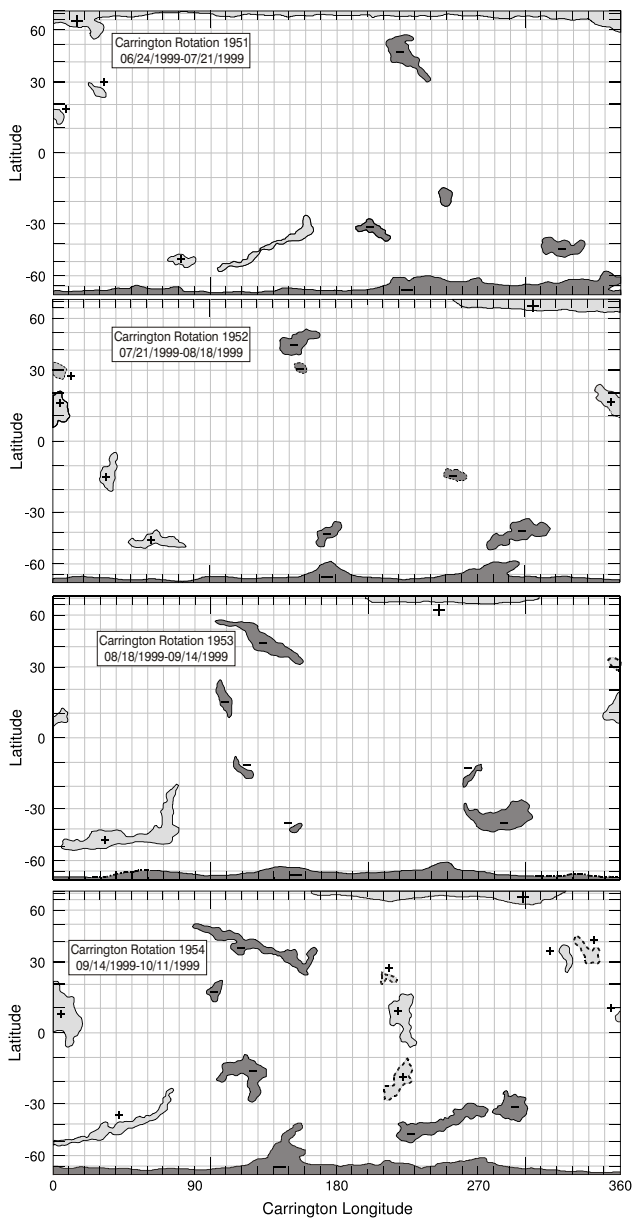


Fig. 5. He I $\lambda 1083$ nm coronal hole maps from Kitt Peak. Light grey is outward polarity and dark grey is inward polarity.

valuable qualitative information about the relative size and distribution of coronal holes at high and low latitudes. At first glance the agreement in Fig. 4 looks fairly good, but the coronal hole maps like those in Fig. 5 show that the latitude extent of the coronal holes is less than the latitude separation between the satellites. In mid 1999 ACE and Ulysses were about 28° apart in latitude, the coronal hole map during this time (Rotation 1951) shows that there are a few small holes. By the end of September (Rotation 1954), many small holes have developed at low latitudes. Furthermore, during this time period coronal holes of like polarity occur near each other, which makes it more difficult to determine if both spacecraft are observing plasma from the same coronal hole.

Examination of ACE and Ulysses superthermal electron data show that counterstreaming occurred at different times at the two spacecraft during the interval shown in Fig. 4, indicating that the structures observed at ACE do not generally map to those observed at Ulysses.

Using the plasma and magnetometer data we find that once the spacecraft are separated by more than 28° , the same structures are not observed at both spacecraft, since after this point the sector structure at both spacecraft starts to become different, but more importantly, the ion and electron measurements indicate different structures. Thus, during the rising phase of 1999 the streams were much smaller than the polar coronal holes observed with Ulysses during solar minimum conditions. However, smaller structures in the heliosphere are consistent with the solar coronal hole maps that show the presence of small coronal holes.

When we examine the data for spacecraft latitude separations between 0° and 42° , there are many instances where both spacecraft appear to observe very similar sectors. We have calculated the difference between the fractions of the inward sector in the model and the measured intervals. Figure 6 shows an example where we have defined boundaries of given regions in the ACE data and then tracked those boundaries by advancing them at each time step in the simulation, using the velocity and acceleration at that grid point. The data are categorized into four groups: high-speed streams (*h*), counterstreaming electrons (*c*), other (*o*), (mostly slow solar wind), and data gaps (*g*). Data gaps were interpolated to allow for long simulation runs. The data gap intervals are not included in our analysis. In Fig. 6 a given region can easily be tracked using the alternating grey and white bands of adjacent regions. For the ACE and Ulysses data curves in Fig. 6, red indicates the outward sector, and black indicates the inward sector. We compared the amount of the inward sector for the Ulysses curve to the amount of the inward sector observed back at ACE for the modeled interval using both the radial propagation and the LARP methods.

In Fig. 7 the percent differences for these two methods of propagation are shown. Points along the 45° line have the same level of agreement for both methods. The source longitude adjustment improves the level of agreement (smaller percent differences) over the radial method. When the source longitude method does not improve the level of agreement, the data points still lie near the 45° line. In Fig. 7, two-thirds of the points lie above the 45° line, which means many intervals are improved with the source longitude adjustment. In Fig. 8 we examine the difference between the data and model fraction of the inward sector for a given interval, as a function of the separation between the spacecraft. Figure 8 shows that the LARP method often improves the sector agreement significantly compared to the RP method. At latitude separations greater than 27° the percentage improvement tends to be larger (black bars). The Ulysses orbit is such that larger latitude separations between the spacecraft occur at smaller radial separations; hence, the percent improvement is larger at distances less than 4 AU.

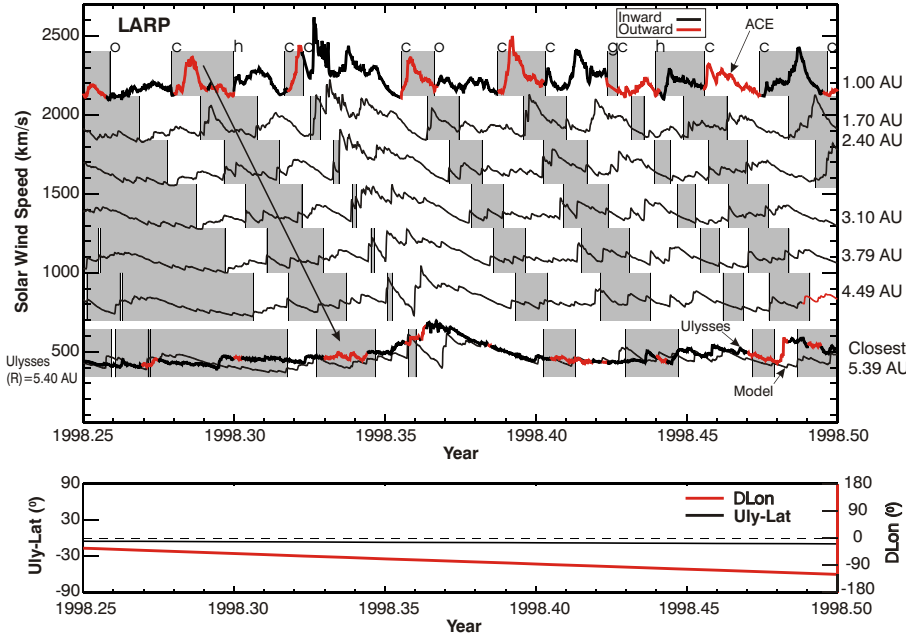


Fig. 6. Same data as in Fig. 1, but now the common source longitude adjustment has been included. The data is divided into regions with vertical bars. These points are then kept track of with in the simulation as they propagate outward. Alternating grey and white shows how regions propagate. The arrow shows how one of the region propagates from 1 AU out 5.4 AU. Above each region is a label indicating the classification. *O* for other, *c* for ICME, *h* for high-speed stream, and *g* for data gap. For the ACE and Ulysses curves, red is outward sector, and black is inward.

5 Conclusions

We utilized a directly driven hydrodynamic model to compare structures observed at ACE in the ecliptic plane near 1 AU with those observed at various heliolatitudes and greater heliocentric distances by Ulysses. The radial variation of the average density and temperature observed in the Ulysses data are consistent with the density being proportional to R^{-2} and temperature proportional to R^{-1} . However, many specific features in the speed, density and temperature are worn down by stream interactions by the time they reach Ulysses.

At smaller latitude separations, particularly less than $\sim 15^\circ$, ACE and Ulysses frequently observe parts of the same stream structures. At greater latitude separations this association tends to break down with little correlation left by the time Ulysses reaches $\sim 30^\circ$ latitude. The association breaks down well before Ulysses reaches high enough southern latitudes, such that only one magnetic polarity is observed. Smith et al. (2001) examined Ulysses observations from years 1999 and 2000, and they found that two magnetic sectors were present up to a latitude of -78° . The LARP method, which corrects for longitude separations between the two spacecraft, generally improves sector alignment, even for greater latitude separations. This reflects the large-scale nature of the sector structure, consistent with the fact that the sector structure generally corotates. The He I $\lambda 1083$ coronal hole maps show that the number of low-latitude coronal holes increases in 1999 and that ACE and Ulysses frequently pass close to, and may be sampling plasma from, different coronal holes once the latitude separation becomes significant. We found many instances where there were small coronal holes of the same magnetic polarity that were near each other and had latitude separations smaller

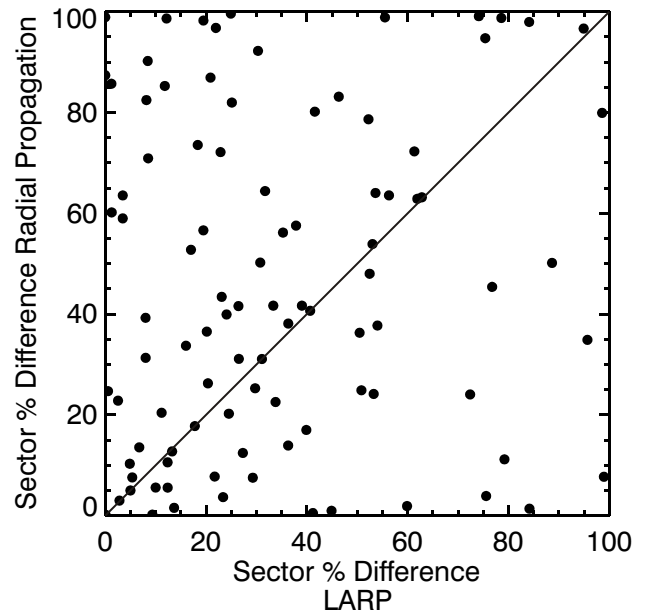


Fig. 7. Percent difference between the amount of inward sector for the model and data intervals. The x-axis uses intervals determined using the radial propagation and the source longitude adjustment. The y-axis is only using the radial propagation.

than the latitude separation of the spacecraft. In such cases the sector structure may be similar, but comparisons between the predicted plasma properties and the Ulysses plasma measurements often break down. Also at these larger latitude separations, instances when the predicted speeds, densities, and temperatures matched frequently appeared to reflect the fact that coronal hole flows tend to have similar plasma prop-

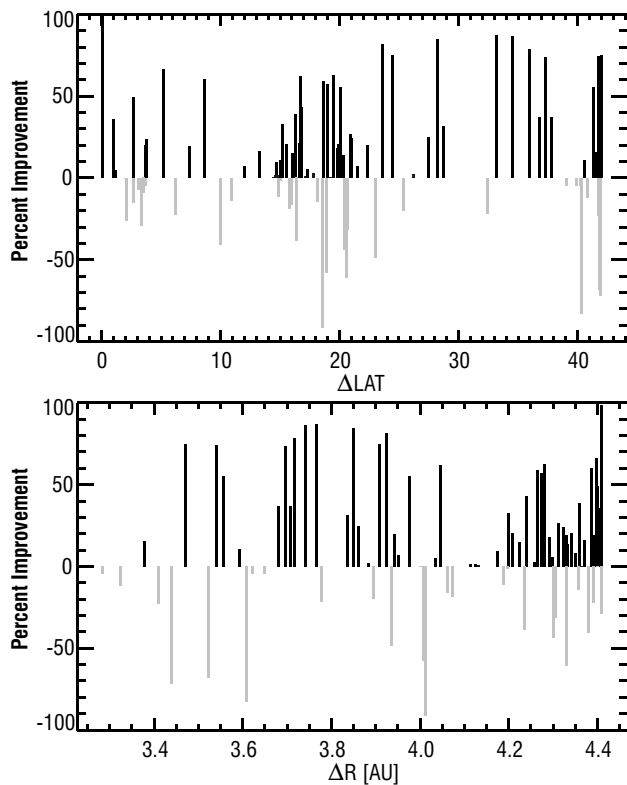


Fig. 8. Amount of improvement in sector agreement. Cases where the source longitude adjustment agreed better with the model are shown in black and cases where the radial propagation along agreed better are shown in grey. The length of the lines represents the differences in the level of agreement for the two methods.

erties. This similarity makes it difficult to determine when ACE and Ulysses were measuring plasma from the same coronal holes. A further complication is that while coronal hole properties are distinct from those of the slow solar wind, flows from Coronal Hole Boundary Layers (CHBLs) provide solar wind plasma with intermediate speeds and heavy ion freezing-in temperatures that span the range between these two types of flows (McComas et al., 2002b); the measured properties of CHBL flows are highly dependent on exactly where a spacecraft passes through them.

In this study we have utilized a directly driven hydrodynamic model to compare observations between widely separated spacecraft. Using ACE and Ulysses measurements we found that during the rising phase of solar cycle 23 stream structures in the heliosphere measured 20–30° in heliolatitude. Our comparison shows that stream structures derived from smaller coronal holes during the rising phase of solar cycle 23 persist over 20–30° in heliolatitude, consistent with solar observations, indicating the presence of small coronal holes.

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